

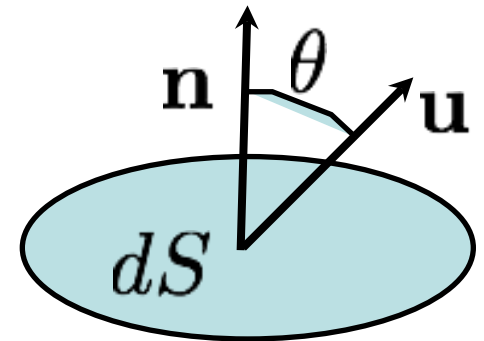
Transfert radiatif entre milieux complexes

Films et métamatériaux

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Rayonnement thermique classique

Luminance $I_\omega(\mathbf{r}, \mathbf{u})$



Flux de chaleur à travers $d\Phi = I_\omega(\mathbf{r}, \mathbf{u}) \cos \theta dS d\Omega$
une surface

Propriétés optiques

- emissivité
- absorptivité

$$I_\omega(\mathbf{r}, \mathbf{u}) = \epsilon_\omega(\mathbf{r}, \mathbf{u}) I_\omega^0(\mathbf{r}, \mathbf{u})$$

$$\epsilon_\omega(\mathbf{r}, \mathbf{u}) = \alpha_\omega(\mathbf{r}, -\mathbf{u})$$

Rayonnement classique: hypothèses et limites

Hypothèses

- Rayon lumineux
- Addition des intensités
- Propriétés locales
- Propriétés surfaciques
(corps opaques)

Limite aux petites échelles

- Diffraction $L \ll \lambda$
- Interférences $L \ll l_{coh}$
- Non localité $L \ll \begin{cases} \lambda \\ l_{coh} \end{cases}$
- Rayonnement volumique $L \ll \delta$

Radiométrie et électromagnétisme

Radiométrie

- Flux d'énergie

$$q = \int I(\mathbf{r}, \mathbf{u}) \mathbf{u} d\Omega$$

- Densité d'énergie

$$u = \int \frac{I(\mathbf{r}, \mathbf{u})}{c} d\Omega$$

- Luminance $I(\mathbf{r}, \mathbf{u})$

Electromagnétisme

- Vecteur de Poynting

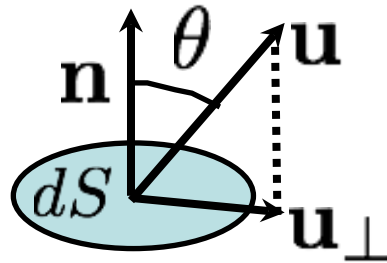
$$\langle \Pi \rangle = \langle \mathbf{E} \times \mathbf{H} \rangle$$

- Densité d'énergie

$$u = \epsilon_0 \frac{E^2}{2} + \mu_0 \frac{H^2}{2}$$

- Luminance ?

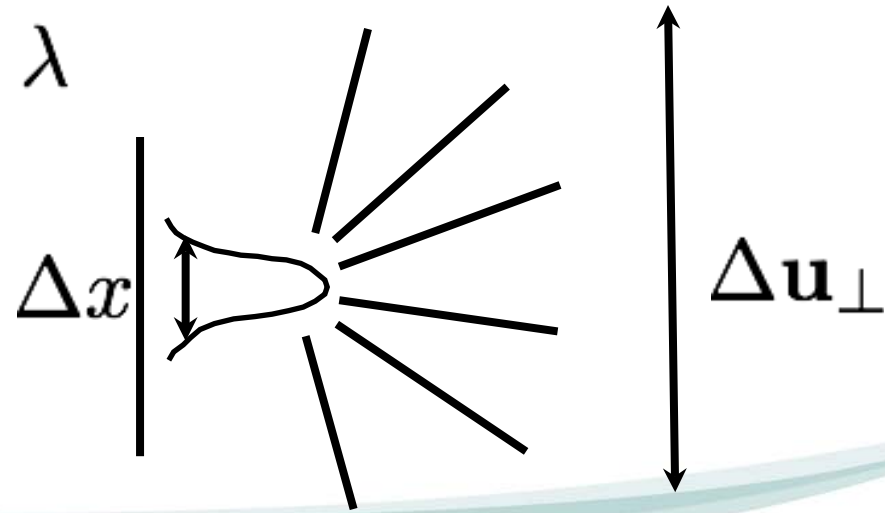
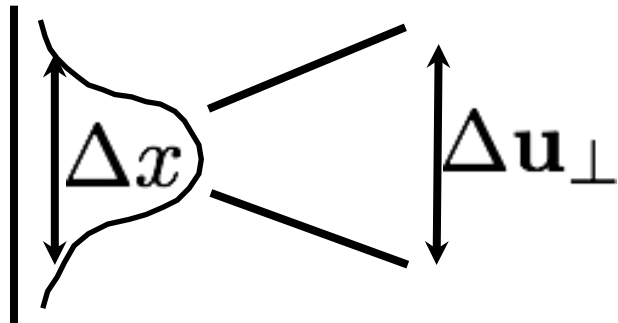
Luminance



Définition statistique de la luminance

$$I(\mathbf{r}, \mathbf{u}) = \left(\frac{k}{2\pi} \right) \cos \theta \int \langle E(\mathbf{r} + \mathbf{r}'/2) E^*(\mathbf{r} - \mathbf{r}'/2) \rangle \exp(-ik\mathbf{u}_\perp \cdot \mathbf{r}') d^3\mathbf{r}'$$

$$\Delta x \Delta \mathbf{u}_\perp \sim \lambda$$

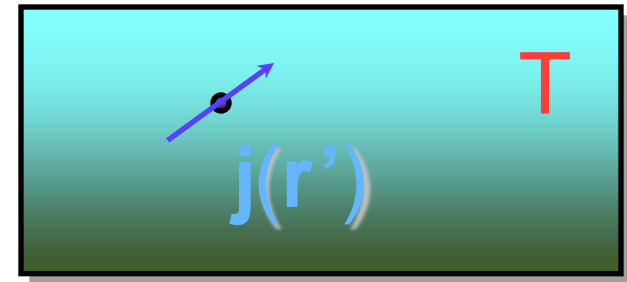


Rayonnement thermique de champ proche

- Electrodynamique fluctuationnelle
- Fonctions de Green
- Théorème de fluctuation-dissipation

$E(\mathbf{r})$

$$\mathbf{E}(\mathbf{r}) = \int \mathbf{G}(\mathbf{r}, \mathbf{r}') \cdot \mathbf{j}(\mathbf{r}') d^3 \mathbf{r}'$$



$$\langle j_n(\mathbf{r}) j_m(\mathbf{r}') \rangle = \frac{\omega \epsilon_0 \text{Im}(\epsilon)}{\pi} \Theta(\omega, T) \delta_{nm} \delta(\mathbf{r} - \mathbf{r}')$$
$$\Theta(\omega, T) = \frac{\hbar \omega}{e^{\hbar \omega / k_b T} - 1}$$

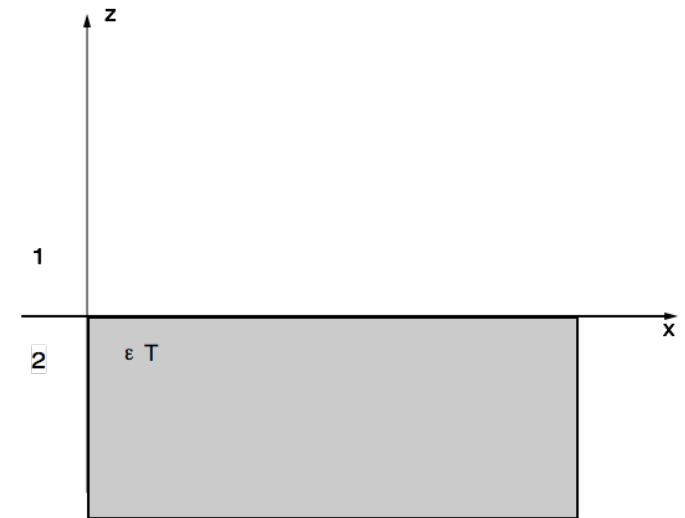
Rayonnement thermique au-dessus d'une surface

Densité d'énergie

$$U(\mathbf{r}, \omega) = \frac{\hbar\omega}{\exp(\hbar\omega/k_bT) - 1} \frac{\omega}{\pi c^2} \text{ImTr} \left[G(\mathbf{r}, \mathbf{r}, \omega) + \frac{c^2 \nabla_{\mathbf{r}}}{\omega^2} \times [\nabla_{\mathbf{r}'} \times^T G(\mathbf{r}, \mathbf{r}', \omega)]^T \Big|_{\mathbf{r}=\mathbf{r}'} \right]$$

Densité d'état

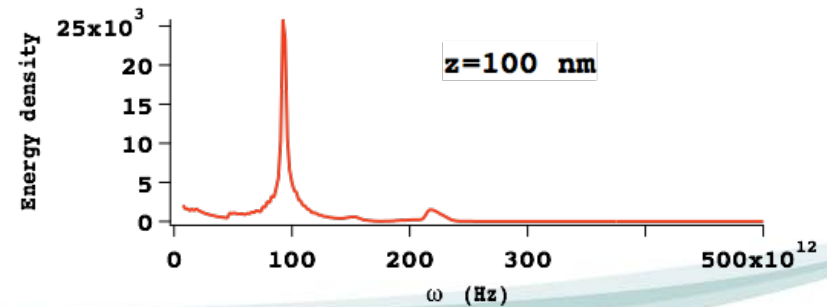
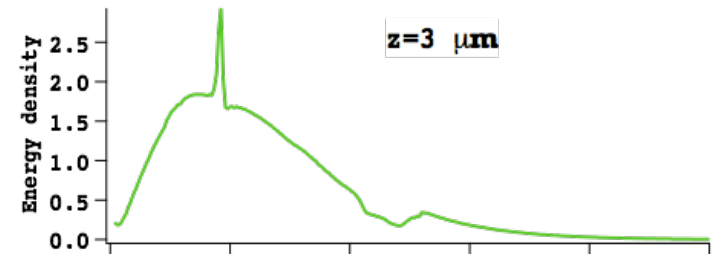
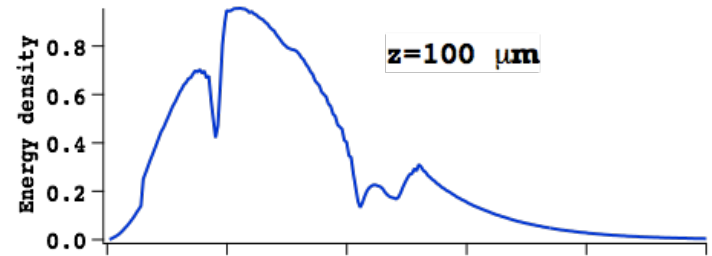
$$U(\mathbf{r}, \omega) = \frac{\hbar\omega}{\exp(\hbar\omega/k_bT) - 1} \rho(\mathbf{r}, \omega)$$



Rayonnement thermique au-dessus d'une surface

- Densité d'énergie

$$U(\mathbf{r}, \omega) = \frac{\hbar\omega^3}{2\pi^2c^3(\exp \hbar\omega/k_bT - 1)} \times \left[\int_0^1 \frac{udu}{\sqrt{1-u^2}} \left\{ 2 + u^2 [Re(r^s e^{2i\sqrt{1-u^2}\omega z/c}) + Re(r^p e^{2i\sqrt{1-u^2}\omega z/c})] \right\} + \int_1^\infty \frac{u^3 du}{\sqrt{u^2-1}} e^{-2\sqrt{1-u^2}\omega z/c} [Im(r^s) + Im(r^p)] \right]$$



Emission thermique au-dessus d'une surface

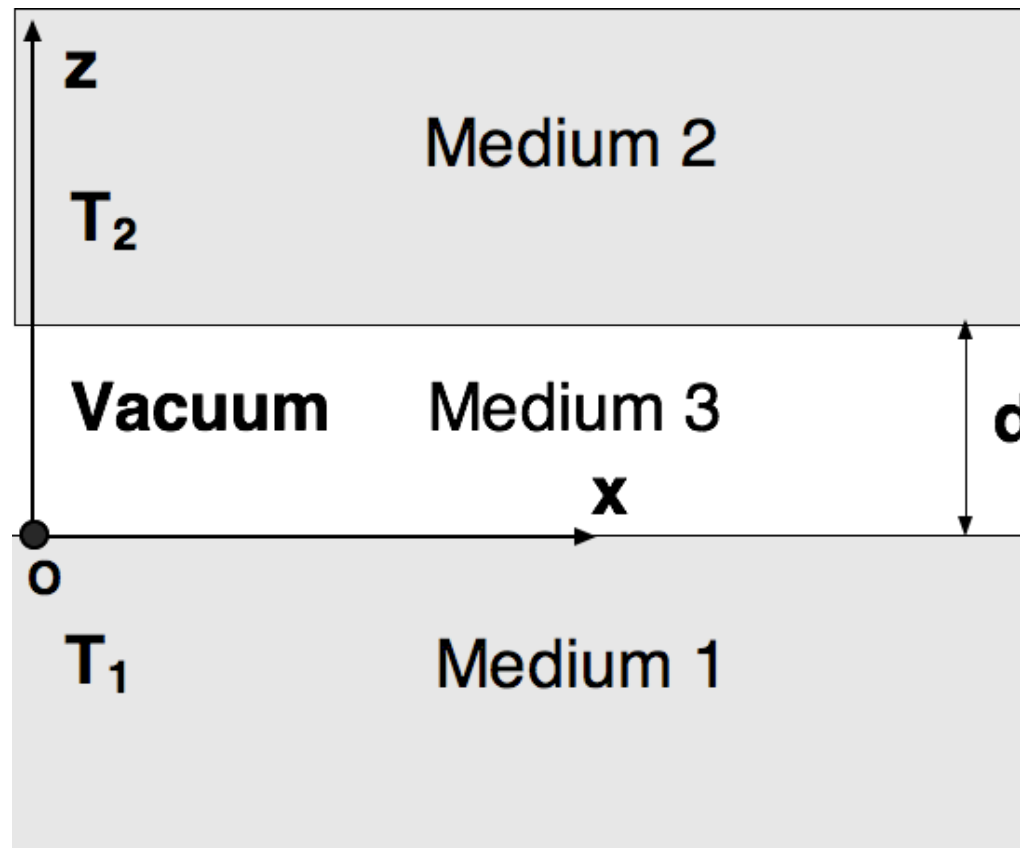
- Densité de flux

$$\langle S_z(\mathbf{r}, \omega) \rangle = \frac{\hbar \omega^3}{2\pi^2 c^2 \exp(\hbar \omega / k_b T) - 1} \times \int_0^1 u du \frac{1 - |r_{12}^s|^2 + 1 - |r_{12}^p|^2}{2}$$

- Emissivité

$$\epsilon = \frac{1 - |r_{12}^s|^2 + 1 - |r_{12}^p|^2}{2}$$

Transfert en champ proche entre deux surfaces planes arbitraires



Transfert en champ proche entre deux surfaces planes arbitraires

- Flux de chaleur monochromatique

$$\langle S_z^{\text{prop}} \rangle = \pi [I_\omega^0(T_1) - I_\omega^0(T_2)] \sum_{\lambda=s,p} \int_0^1 u du \frac{(1 - |r_{31}^\lambda|^2)(1 - |r_{32}^\lambda|^2)}{|1 - r_{31}^\lambda r_{32}^\lambda e^{2i\sqrt{1-u^2}\omega d/c}|^2}$$

$$\langle S_z^{\text{evan}} \rangle = 4\pi [I_\omega^0(T_1) - I_\omega^0(T_2)] \sum_{\lambda=s,p} \int_1^\infty u du \frac{\text{Im}(r_{31}^\lambda) \text{Im}(r_{32}^\lambda) e^{-2\sqrt{u^2-1}\omega d/c}}{|1 - r_{31}^\lambda r_{32}^\lambda e^{-2\sqrt{u^2-1}\omega d/c}|^2}$$

G. Bimonte , Phys. Rev. A, **80**, 042102 (2009)

Interprétation en termes de nombre de modes couplés

- Formalisme de Landauer

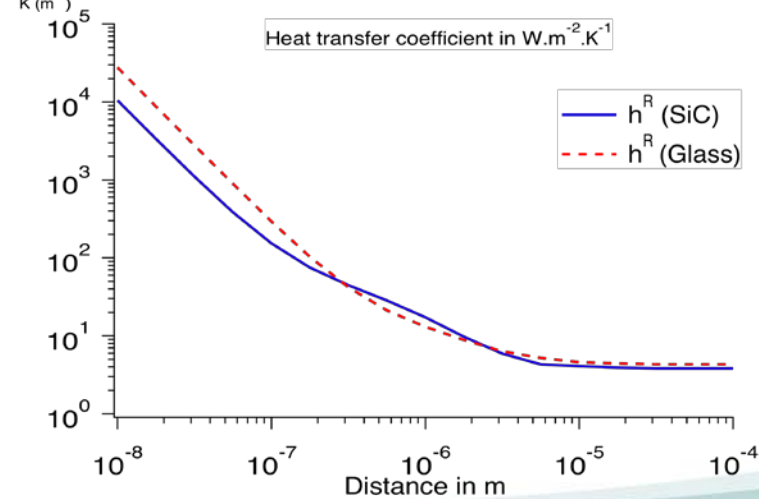
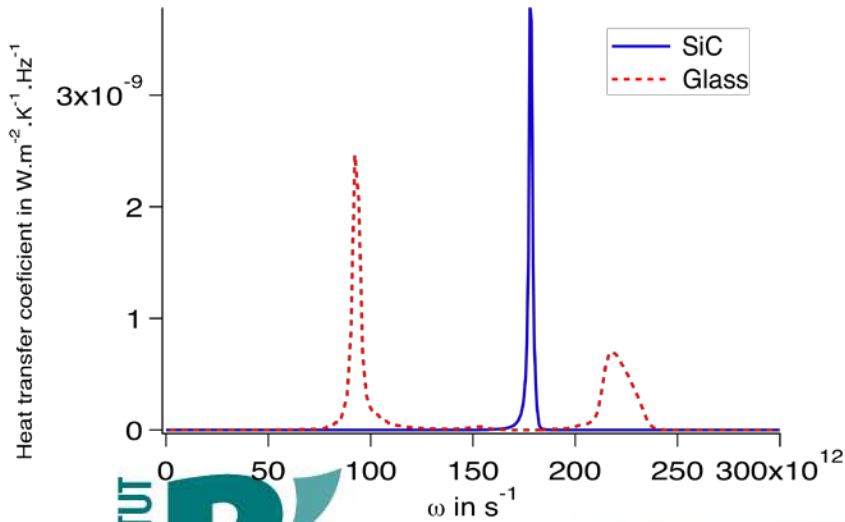
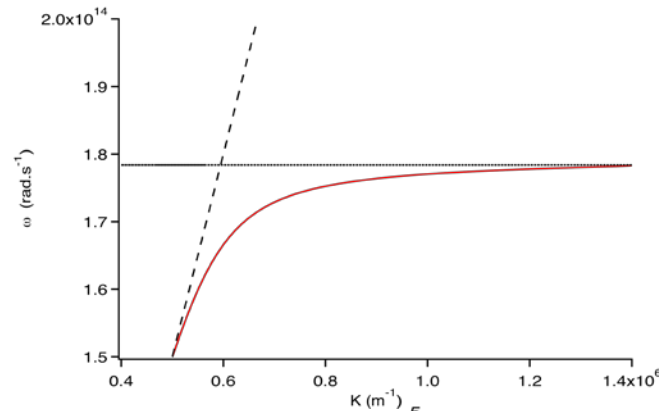
$$\varphi = \int_0^{\infty} \frac{\hbar\omega}{\exp(\hbar\omega/k_bT) - 1} N(\omega) \frac{d\omega}{2\pi}$$

- Nombre de modes couplés

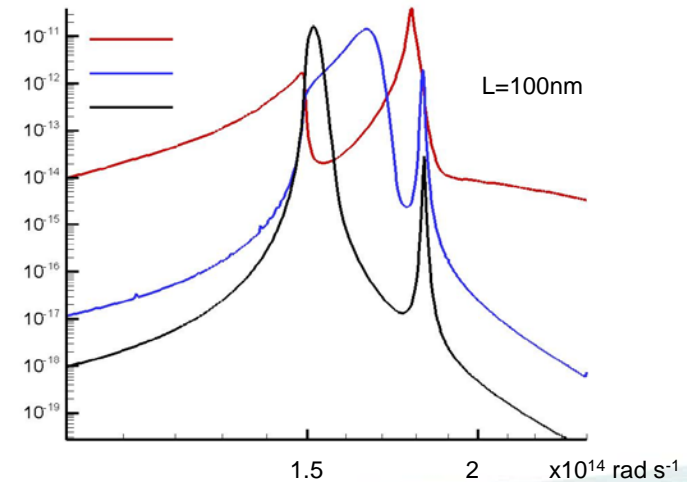
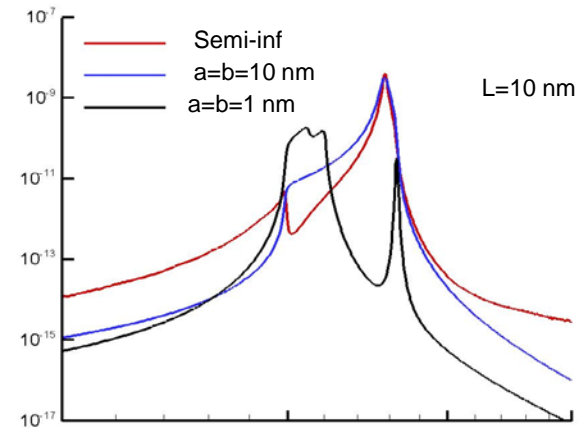
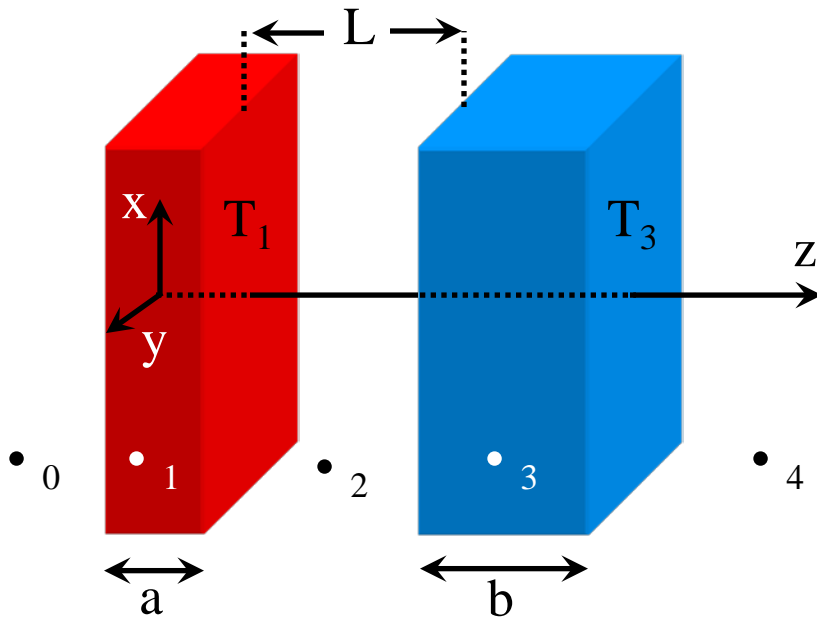
$$N(\omega) = \langle S_z(\omega) \rangle \times \frac{2\pi(\exp(\hbar\omega/k_bT) - 1)}{\hbar\omega}$$

Augmenter le transfert thermique

Surface polariton

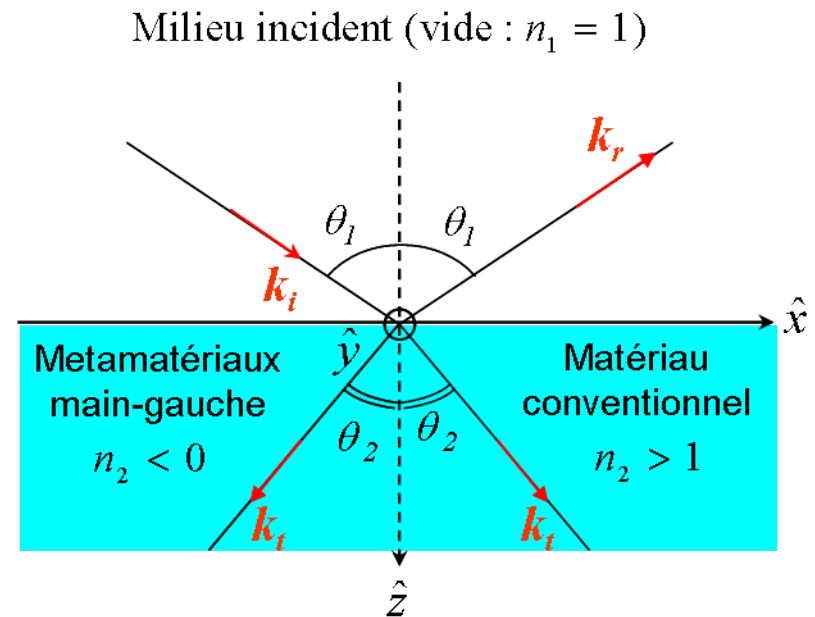
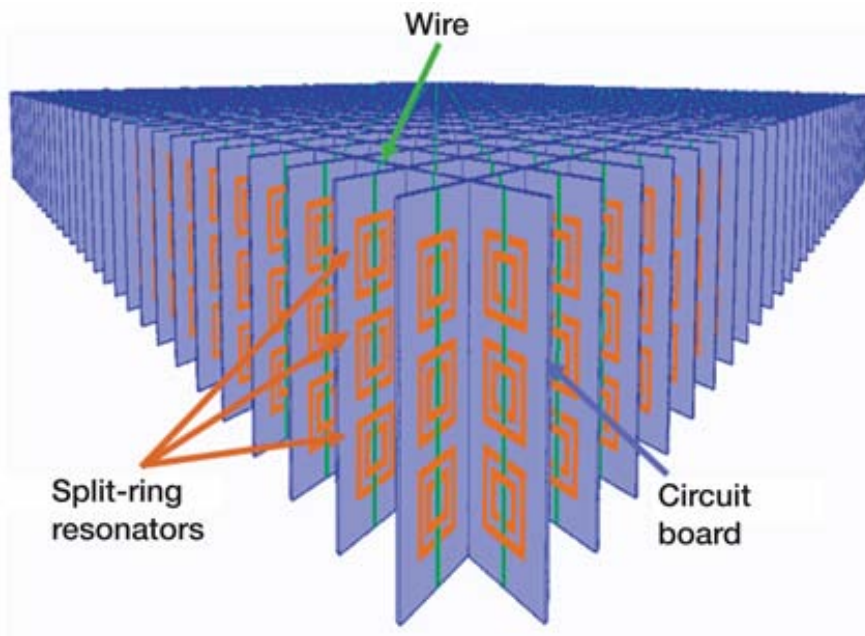


Transfert entre films



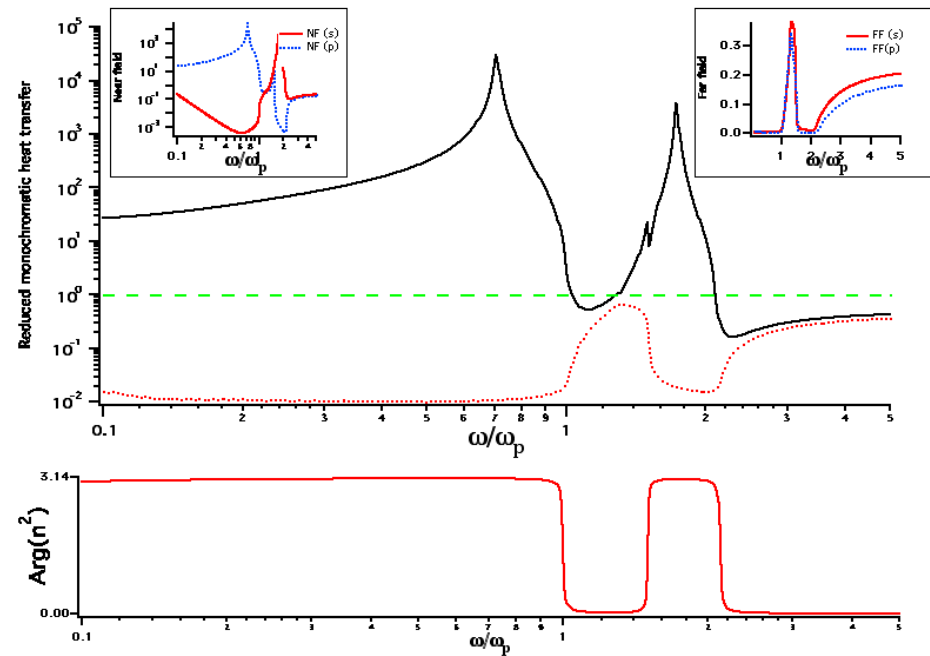
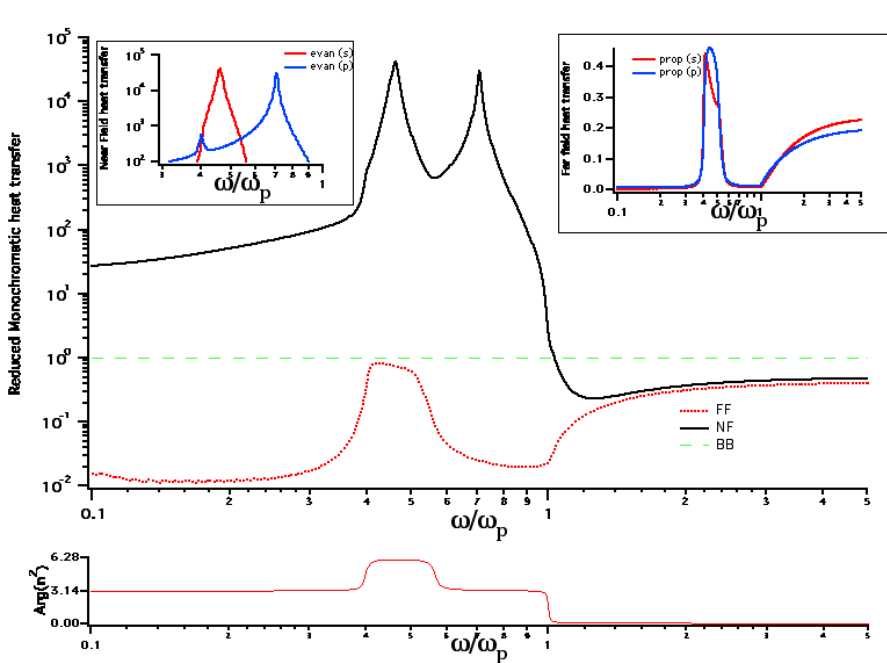
Transfert entre métamatériaux

- Métamatériaux

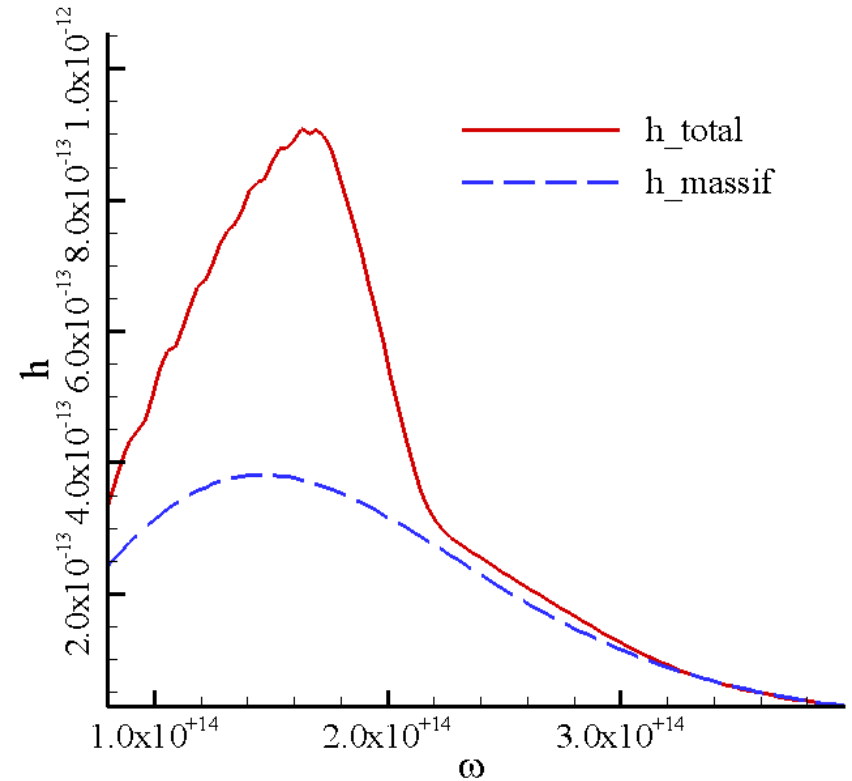
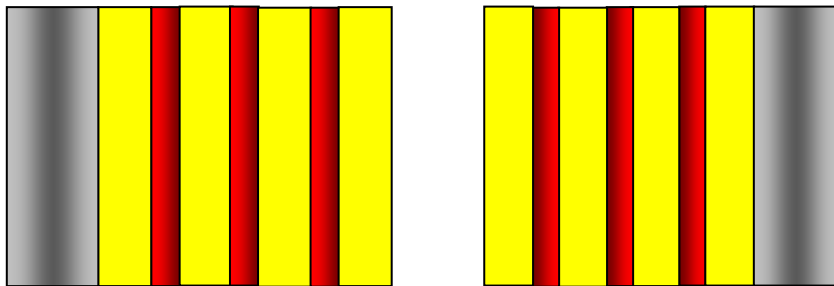


Transfert entre métamatériaux

- Métamatériaux



Transfert entre cristaux photoniques



Perspectives

- Trouver les limites du transfert à petite échelles
- Augmenter le nombre de modes couplés entre corps chauffés
- Applications : augmenter le niveau de signal TPV.